ACTA TECHNICA CORVINIENSIS
– Bulletin of Engineering
Tome IX [2016] , Fascicule 1 [January – March]
ISSN: 2067 – 3809

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LAPLACE TRANSFORM SOLUTIONS FOR MAGNETO-HYDRODYNAMIC BOUNDARY LAYER FLOW AND HEAT TRANSFER IN A POROUS MEDIUM WITH THERMAL RADIATION EFFECT

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Abstract: A mathematical model is developed for unsteady Magnetohydrodynamic boundary layer flow and heat transfer through a Darcian porous medium bounded by a uniformly moving semi-infinite isothermal vertical plate in presence of thermal radiation. The flow model is considered as an viscous, incompressible, electrically-conducting Newtonian fluid which is an optically thin gray gas. Suitable transformations are used to convert the partial differential equations corresponding to the momentum and energy equations into nonlinear ordinary differential equations. Analytical solutions of these equations are obtained by Laplace transform. The effects of Hartmann number (M), porosity parameter (K), thermal radiation parameter (R_a), and Prandtl number (Pr) on flow velocity, fluid temperature, velocity and temperature gradients at the surface are studied graphically. Velocity is reduced with Hartmann number but enhanced with thermal radiation and porosity parameter. An increase in porosity/thermal radiation parameter is found to strongly enhance flow velocity values. Velocity gradient at y=0 is increased with porosity parameter. Applications of the study arise in engineering and geophysical sciences like magnetohydrodynamic transport phenomena and magnetic field control of materials processing, solar energy collector systems.
Keywords: optically thin gray gas; Hartmann number; porous media; heat transport; unsteady boundary layer flow

INTRODUCTION

Fluid flow through a porous media has been studied theoretically and experimentally by numerous authors due to its wide applications in various fields such as diffusion technology, transpiration cooling, hemodialysis processes, flow control in nuclear reactors, etc. In view of geophysical applications of the flow through porous medium, a series of investigations has been made by Raptis et.al (1981-1982), where the porous medium is either bounded by horizontal, vertical surfaces or parallel porous plates. Singh et.al (1989) and Lai and Kulacki (1990) have been studied the free convective flow past vertical wall. Nield (1994) studied convection flow through porous medium with inclined temperature gradient. Singh et al. (2005) also studied periodic solution on oscillatory flow through channel in rotating porous medium. Further due to increasing scientific and technical applications on the effect of radiation on flow characteristic has more importance in many engineering processes occurs at very high temperature and acknowledge radiative heat transfer such as nuclear power plant, gas turbine and various propulsion devices for aircraft, missile and space vehicles. The effect of radiation on flow past different geometry a series of investigation have been made by Hassan (2003), Seddeek (2000) and Sharma et al (2011). The combined radiation–convection flows have been extended by by Ghosh and Be'g (2008) to unsteady convection in porous media. Hossain and Takhar (1996) studied the mixed convective flat plate boundary-layer problem using the Rosseland (diffusion) flux model. Mohammadein et al. (1998) studied the radiative flux effects on free convection in the Darcian porous media using the Rosseland model. The transient magnetohydrodynamic free convective flow of a viscous, incompressible, electrically conducting, gray, absorbing–emitting, but non-scattering, optically thick fluid medium which occupies a semi-infinite porous region adjacent to an infinite hot vertical plate moving with a constant velocity is presented by Ahmed and Kalita (2013). Raptis and Perdikis (2004) have also studied analytically the transient convection in a highly porous medium with unidirectional radiative flux. Ghosh and Pop (2007) studied indirect radiation effects on convective gas flow. Ahmed and Kalita (2013) investigated the effects of chemical reaction as well as magnetic field on the heat and mass transfer of Newtonian two-dimensional flow over an infinite vertical oscillating plate with variable mass diffusion. Ahmed (2014) presented the effects of conduction-radiation, porosity and chemical reaction on unsteady hydromagnetic free convection flow past an impulsively-started semi-infinite vertical plate embedded in a porous medium in presence of thermal radiation. The thermal radiation and Darcian drag force MHD unsteady thermal-convection flow past a semi-infinite vertical plate immersed in a semi-infinite saturated porous regime with variable surface temperature in the presence of transversal uniform magnetic field have been discussed by Ahmed el al. (2014).

The present paper is to investigate the effect of magnetic field and radiation on unsteady free convection heat transfer flow of viscous laminar electrically conducting Newtonian radiating fluid past an impulsively started semi-infinite vertical surface in a Darcian porous medium. The analytical solution is obtained using Laplace Transform technique and discussed graphically for various flow parameters.

MATHEMATICAL FORMULATION

Considering the magneto-hydrodynamic unsteady free convection and heat transfer flow of a viscous, incompressible, electrically conducting Newtonian fluid past a semi-infinite isothermal vertical plate embedded in a porous media under the influence of the thermal buoyancy.

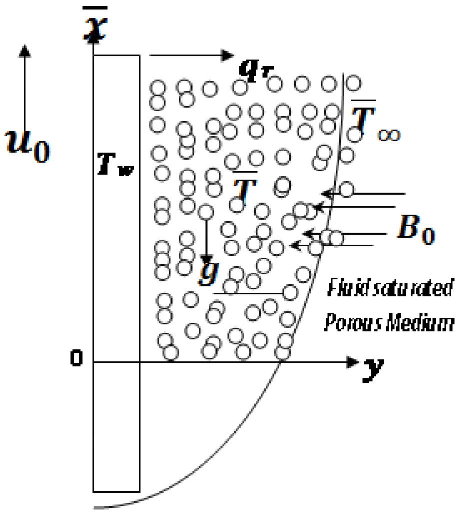

Figure 1: Physical model and coordinate system

A uniform magnetic filed of uniform strength B_0^2 is assumed to be applied normal to the surface. The flow is assumed to be in the $\bar{x}$-direction, which is taken along the plate in the upward direction and $\bar{y}$-axis is normal to it. Initially it is assumed that the plate and the fluid are at the same temperature $\bar{T}$. At time t>0, the plate temperature is instantly raised to $\bar{T}_w > \bar{T}_\infty$ and, which is thereafter maintained constant, where $\bar{T}_\infty$ is the temperature outside the boundary layer. The induced magnetic field and viscous dissipation is assumed to be negligible as the magnetic Reynolds number of the flow is taken to be very small. Assuming that the Boussinesq and boundary-layer approximations hold, the governing equations to the problem are given by:

$$\frac{\partial \bar{u}}{\partial \bar{t}} = g\beta(\bar{T}-\bar{T}_\infty) + \nu\frac{\partial^2 \bar{u}}{\partial^2 \bar{y}} - \frac{\sigma B_0^2}{\rho}\bar{u} - \frac{\nu}{\bar{K}}\bar{u} , \qquad (1)$$

$$\rho C_p \frac{\partial \bar{T}}{\partial \bar{t}} = \kappa\frac{\partial^2 \bar{T}}{\partial \bar{y}^2} - \frac{\partial q_r}{\partial \bar{y}} . \qquad (2)$$

The initial and boundary conditions are

$$\left.\begin{aligned} &\bar{u}=0, \bar{T}=\bar{T}_\infty, \forall \bar{y}, \bar{t}\le 0 \\ &\bar{u}=u_0, \bar{T}=\bar{T}_w \text{ at } \bar{y}=0, \bar{t}>0 \\ &\bar{u}=0, \bar{T}=\bar{T}_\infty, \text{ as } \bar{y}\to\infty, \bar{t}>0 \end{aligned}\right\} \qquad (3)$$

The local radiant absorption for the case of an optically thin gray gas is expressed (Cogley et al. (1968)) as

$$\frac{\partial q_r}{\partial \bar{y}} = -4\bar{a}\bar{\sigma}\left(\bar{T}_\infty^4 - \bar{T}^4\right) , \qquad (4)$$

where $\bar{\sigma}$ and $\bar{a}$ are the Stefan-Boltzmann constant and mean absorption co-efficient respectively. We assume that the differences within the flow are sufficiently small so that $\bar{T}^4$ can be expressed as a linear function of $\bar{T}$ after using Taylor's series to expand $\bar{T}^4$ about the free stream temperature $\bar{T}_\infty^4$ and neglecting higher order terms. This results in the following approximation:

$$\bar{T}^4 \cong 4\bar{T}_\infty^3\bar{T} - 3\bar{T}_\infty^4 , \qquad (5)$$

$$\rho C_p \frac{\partial \bar{T}}{\partial \bar{t}} = \kappa\frac{\partial^2 \bar{T}}{\partial \bar{y}^2} - 16\bar{a}\bar{\sigma}\bar{T}_\infty^3\left(\bar{T}-\bar{T}_\infty\right) . \qquad (6)$$

Introducing the following non-dimensional quantities:

$$y=\frac{\bar{y}u_0\sqrt{G}}{\nu} , \; u=\frac{\bar{u}}{u_0} , \; M=\frac{\sigma B_0^2\nu}{\rho u_0^2 G} , \; K=\frac{u_0^2\bar{K}G}{\nu^2} ,$$

$$G=\frac{g\beta\nu\left(\bar{T}_w-\bar{T}_\infty\right)}{u_0^3} , \quad \theta=\frac{\bar{T}-\bar{T}_\infty}{\bar{T}_w-\bar{T}_\infty} , \; Pr=\frac{\mu C_p}{\kappa}$$

$$t=\frac{\bar{t}u_0^2G}{\nu} , \; R_a=\frac{16\bar{a}\bar{\sigma}\nu^2\bar{T}_\infty^3}{\kappa u_0^2} , \; \nu=\frac{\mu}{\rho} . \qquad (7)$$

Using the transformations (7), the non-dimensional forms (1), (3) and (6) are

$$\frac{\partial u}{\partial t} = Gr\theta + \frac{\partial^2 u}{\partial y^2} - (M+K^{-1})u , \qquad (8)$$

$$\frac{\partial\theta}{\partial t}=\frac{1}{Pr}\frac{\partial^2\theta}{\partial y^2}-\frac{R_a}{Pr}\theta\ . \quad (9)$$

The corresponding initial and boundary conditions transformed to:

$$\left.\begin{array}{l} u=0,\theta=0,\forall y,t\le 0 \\ u=1,\theta=1 \text{ at } y=0,t>0 \\ u=0,\theta=0, \text{as } y\to\infty,t>0 \end{array}\right\} \quad (10)$$

METHOD OF SOLUTION

The unsteady, non-linear, coupled partial differential equations (8) and (9) along with their boundary conditions (10) have been solved analytically using Laplace transforms technique and their solutions are as follows:

$$u(y,t)=\frac{1}{2}\left[\begin{array}{l}\left(1-\frac{1}{\psi}\right)\left\{e^{2\eta\sqrt{\xi t}}\operatorname{erfc}\left(\eta+\sqrt{\xi t}\right)+e^{-2\eta\sqrt{\xi t}}\operatorname{erfc}\left(\eta-\sqrt{\xi t}\right)\right\} \\ +\frac{1}{\psi}e^{\lambda t}\left\{\begin{array}{l}e^{2\eta\sqrt{(\xi+\lambda)t}}\operatorname{erfc}\left(\eta+\sqrt{(\xi+\lambda)t}\right) \\ +e^{-2\eta\sqrt{(\xi+\lambda)t}}\operatorname{erfc}\left(\eta-\sqrt{(\xi+\lambda)t}\right)\end{array}\right\} \\ +\frac{1}{\psi}\left\{\begin{array}{l}e^{2\eta\sqrt{R_a t}}\operatorname{erfc}\left(\eta\sqrt{Pr}+\sqrt{R_a t}\right) \\ +e^{-2\eta\sqrt{R_a t}}\operatorname{erfc}\left(\eta\sqrt{Pr}-\sqrt{R_a t}\right)\end{array}\right\} \\ -\frac{1}{\psi}e^{\lambda t}\left\{\begin{array}{l}e^{2\eta\sqrt{(R_a+\lambda)t}}\operatorname{erfc}\left(\eta\sqrt{Pr}+\sqrt{(R_a+\lambda)t}\right) \\ +e^{-2\eta\sqrt{(R_a+\lambda)t}}\operatorname{erfc}\left(\eta\sqrt{Pr}-\sqrt{(R_a+\lambda)t}\right)\end{array}\right\}\end{array}\right], \quad (11)$$

$$\theta(y,t)=\frac{1}{2}\left\{\begin{array}{l}e^{2\eta\sqrt{R_a t}}\operatorname{erfc}\left(\eta\sqrt{Pr}+\sqrt{R_a t}\right) \\ +e^{-2\eta\sqrt{R_a t}}\operatorname{erfc}\left(\eta\sqrt{Pr}-\sqrt{R_a t}\right)\end{array}\right\}, \quad (12)$$

where $\xi=M+K^{-1}$, $\eta=\frac{y}{2\sqrt{t}}$, $\psi=\xi-R_a$, $\lambda=\frac{\psi}{Pr-1}$.

SKIN FRICTION AND NUSSELT NUMBER

The non-dimensional skin friction and Nusselt number is given as follows:

$$\tau=-\left[\frac{\partial u(y,t)}{\partial y}\right]_{y=0}$$

$$\left.\begin{array}{l}=\left(1-\frac{1}{\psi}\right)\left\{\frac{e^{-\xi t}}{\sqrt{\pi t}}+\sqrt{\xi}\operatorname{erf}\left(\sqrt{\xi t}\right)\right\} \\ +\frac{1}{\psi}e^{\lambda t}\left\{\frac{e^{-(\xi+\lambda)t}}{\sqrt{\pi t}}+\sqrt{(\xi+\lambda)}\operatorname{erf}\left(\sqrt{(\xi+\lambda)t}\right)\right\} \\ +\frac{1}{\psi}\sqrt{Pr}\left\{\frac{e^{-R_a t}}{\sqrt{\pi t}}+\sqrt{R_a}\operatorname{erf}\left(\sqrt{R_a t}\right)\right\} \\ -\frac{1}{\psi}\sqrt{Pr}e^{\lambda t}\left\{\frac{e^{-(R_a+\lambda)t}}{\sqrt{\pi t}}+\sqrt{(R_a+\lambda)}\operatorname{erf}\left(\sqrt{(R_a+\lambda)t}\right)\right\}\end{array}\right\} \quad (13)$$

$$Nu=-\left[\frac{\partial\theta(y,t)}{\partial y}\right]_{y=0}=\sqrt{Pr}\left\{\frac{e^{-R_a t}}{\sqrt{\pi t}}+\sqrt{R_a}\operatorname{erf}\left(\sqrt{R_a t}\right)\right\}. \quad (14)$$

RESULTS AND DISCUSSION

The problem of thermal radiation effect on a porous media transport under optically thick approximation formulated, analyzed and solved analytically. In order to point out the effects of physical parameters namely; magnetohydrodynamic force (M), radiation parameter (R_a), Porosity parameter (K) on the flow patterns, the computation of the flow fields are carried out. The values of velocity, temperature, shear stress and rate of heat transfer are obtained for the physical parameters as mention. The velocity profiles has been studied and presented in Figures 2 to 4. Figure 2 shows the effect of the Hartmann number M on the fluid velocity and the results show that the presence of the magnetic force causes retardation of the fluid motion represented by general decreases in the fluid velocity. It is due the fact that magnetic force which is applied in the normal direction to the flow produces a drag force which is known as Lorentz force.

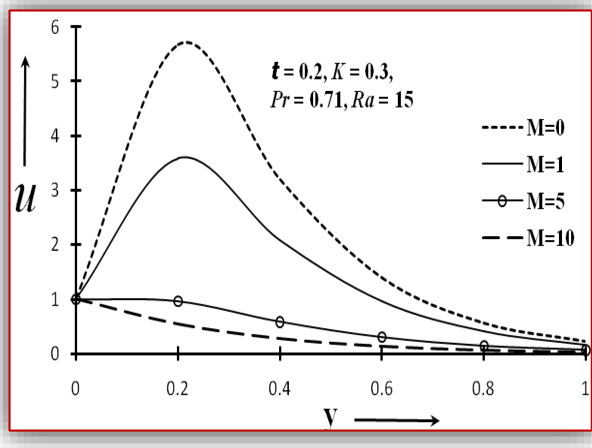

Figure 2: Flow velocity distribution for Hartmann number M

The opposite trend is observed in Figure 3 for the case when the value of the porous permeability (K = 0.2, 0.5, 1.0, 1.5) is increased. As depicted in this figure, the effect of increasing the value of porous permeability is to increase the value of the velocity component in the boundary layer due to the fact that drag is reduced by increasing the value of the porous permeability on the fluid flow which results in increased velocity. The trend shows that the velocity is accelerated with increasing porosity parameter. The effect of velocity for different values of radiation (R_a =0, 15, 16, 18) is also presented in Figure 4. It is then observed that the flow velocity is accelerated with increasing values of radiation. Also it is seen that without radiation (R_a =0, Figure 4) or for the small value K = 0.2 (Figure 3), the values of flow velocity reduces exponentially from the plate, while for the higher values of K or R_a the flow velocity has a bigger pick in the neighbourhood of y = 0.2, but the opposite behaviour has been observed for the effects of higher magnetohydrodynamic force (M=10, Figure 2).

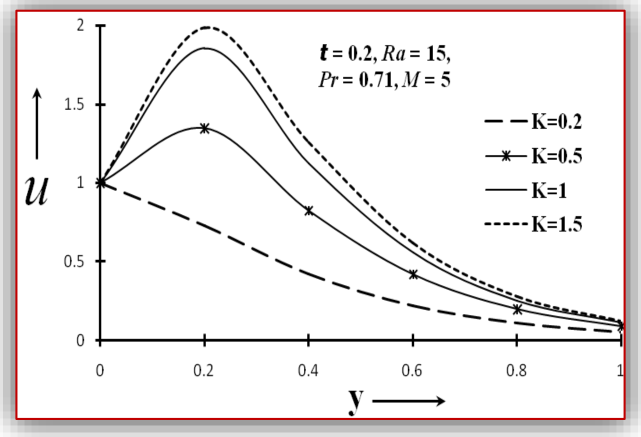

Figure 3: Flow velocity distribution for porosity K

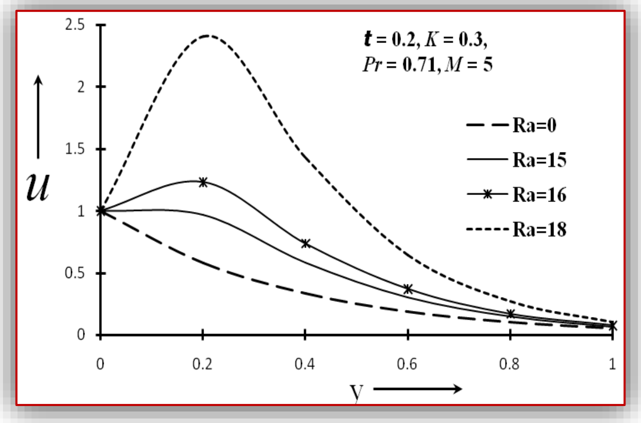

Figure 4: Flow velocity distribution for radiation R_a

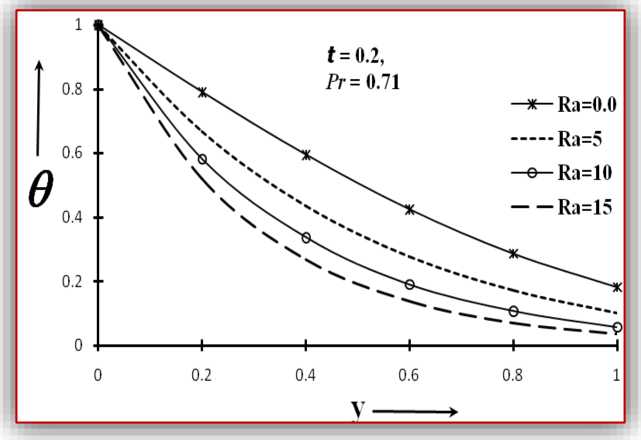

Figure 5: Temperature distribution for radiation R_a

The temperature profiles are calculated for different values of thermal radiation parameter (R_a=0, 5, 10, 15) at time t = 0.2 and these are shown in Figure 5. The effect of thermal radiation parameter is important in temperature profiles. It is observed that the temperature increases with decreasing radiation parameter. Figure 6 reveals temperature variations with Pr (Prandtl number) which signifies the ratio of momentum to thermal diffusivity at t = 0.2. The temperature is observed to decrease with an increase in Pr. For lower Pr fluids, heat diffuses faster than momentum and vice versa for higher Pr fluids. Larger Pr values correspond to a thinner thermal boundary layer thickness and more uniform temperature distributions across the boundary layer. Smaller Pr fluids possess higher thermal conductivities so that heat can diffuse away from the vertical surface faster than for higher Pr fluids (low Pr fluids correspond to thicker boundary layers). For working oils (Pr = 11.4), convection is very effective in transferring energy from an area, compared to pure conduction and momentum diffusivity is dominant. It is also observed that the temperature is maximum near the plate and decreases away from the plate and finally takes asymptotic value for all values of Pr.

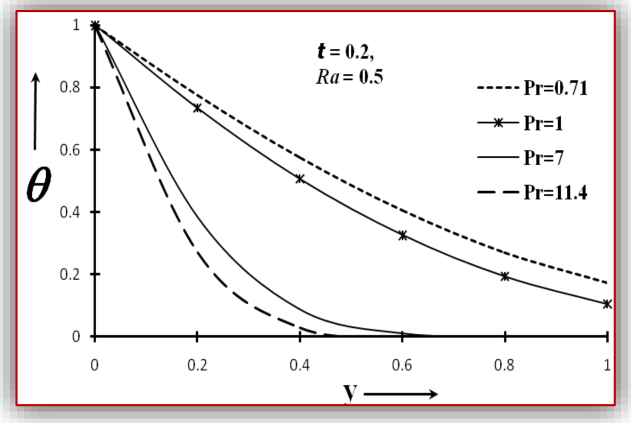

Figure 6: Temperature distribution for Pr

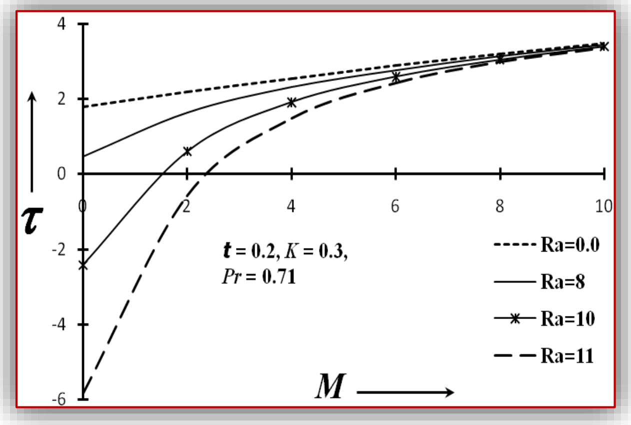

Figure 7: Skin friction distribution for radiation R_a

Figure 7 illustrates the transient shear stress variation with Hartmann number and radiation parameter. The shear stresses at the wall are seemed to be enhanced with a rise in Hartmann number, which is proportional to the square of the magnetic field, B_0. A reversed trend has been observed for conduction-radiation on shear stress (τ) i.e. τ decreases substantially at the wall for R_a = 0, 8, 10, 11. For the non-radiating flow case, R_a = 0, a significant linear flow of shear stress is sustained against hydromagnetic force. For the case, R_a = 10, 11, a significant flow reversal (backflow) is obtained within the region 0<M<2.5 i.e. shear stresses become negative. However for R_a = 0, 8, all backflow is eliminated entirely from the regime for all hydromagnetic forces and only positive shear stresses arise at the plate.

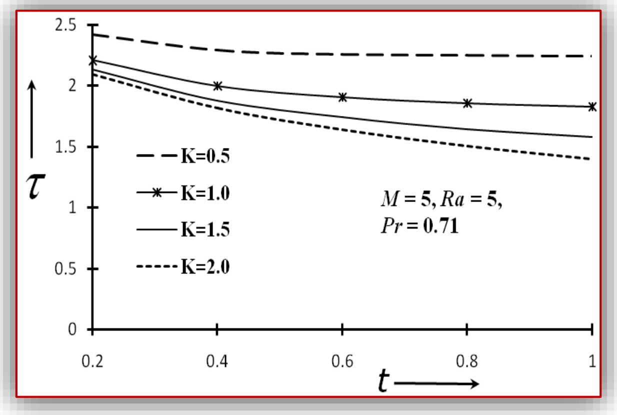

Figure 8: Skin friction distribution for Pr

Figure 8 shows the distribution of shear stress at the wall for various porosity parameters over time. With a rise in radiation parameter, K, from 0.5, 1.0 through 1.5 to 2.0 decreases the magnitude of the shear stress through the boundary layer. We observe that for all values of K, shear stress remains positive i.e. no flux reversal arises for all times into the boundary layer. With progression in time, t, the shear stress is however found to decrease continuously. Finally, in Figure 9 the distribution of rate of heat transfer with radiation parameter is shown against t. Inspection shows that, increasing radiation parameter, R_a, tends to boost the heat transfer rate at the wall i.e. elevate Nu magnitudes. A substantial decrease is observed in Nu for the time parameter.

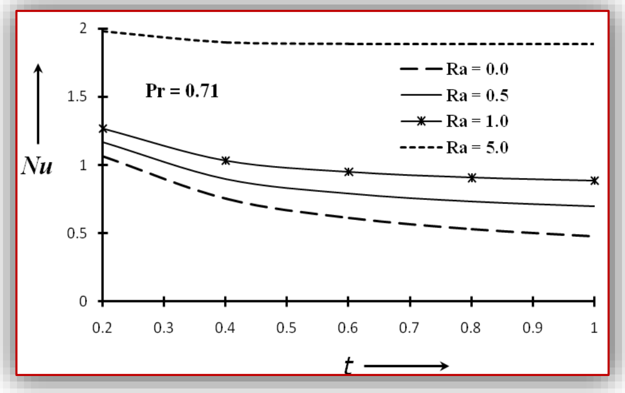

Figure 9: Nusselt number distribution for radiation R_a

CONCLUSION

In the present work, we have analyzed flow, heat transfer on convection flow of a viscous incompressible, electrically conducting and radiating fluid over an infinite vertical plate embedded in a Darcian porous regime in the presence of transverse magnetic field and thermal radiation using the classical model for the radiative heat flux. Final results are computed for variety of physical parameters which are presented by means of graphs. Laplace transforms solutions for the non-dimensional momentum and energy equations subject to transformed boundary conditions have been obtained and the results indicate that:

≡ The flow has been shown to be decelerated with increasing Hartmann number but accelerated with conduction-radiation and porosity parameters.
≡ Increasing Hartmann number also increases the shear stress and back flow has been observed for higher radiation near the wall.
≡ A positive decrease in R_a or K strongly enhanced the shear stress.
≡ Increasing thermal radiation contribution (R_a) serves to enhance wall heat transfer gradient significantly in the porous regime.
≡ With an increase in time (t), both the skin friction and wall heat transfer are decreased.
≡ Temperature is decreased with an increase in thermal radiation contribution (R_a).

The study has important applications in materials processing and nuclear heat transfer control, as well as MHD energy generators. The current study has employed a Newtonian viscous model. Presently the authors are extending this work to consider viscoelastic fluids and also power-law rheological fluids. The results of these studies will be presented imminently.

NOMENCLATURE

u non- dimensional velocity component in x direction [ms^{-1}]
y normal direction of vertical plane surface [m]
C_p specific heat at constant pressure [J $Kg^{-1}K^{-1}$]
D chemical molecular diffusivity [m^2s^{-1}]
G acceleration due to gravity [ms^{-2}]
G free convection parameter [-]
M Hartmann number (magnetic parameter) [-]
K permeability of the porous medium [m^2]
Pr Prandlt number [-]
P pressure [mmHg]
θ temperature [K]
$\overline{T}$ dimensional temperature
$\overline{T}_w$ dimensional temperature at the plate
$\overline{T}_\infty$ dimensional temperature at the free stream
T non-dimensional time [S]
u_0 plate velocity
B_0 strength of the magnetic field
R_a Radiation parameter
q_r Radiative heat flux

$\bar{\sigma}$ Stefan-Boltzmann constant
$\bar{a}$ mean absorption co-efficient

Greek symbols

β volumetric coefficient of thermal expansion [K^{-1}]
κ thermal conductivity, [$J.m^{-1}s^{-1}K^{-1}$]
μ kinematic viscosity [m^2s^{-1}]
ρ density [Kgm^{-3}]
σ electrical conductivity
τ coefficient of viscosity

Subscripts

w conditions on the plane surface
∞ conditions away from the plane surface

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